

Multiple zeta values in deformation quantization

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Hamiltonian mechanics

Particle in 1d:

position

x

momentum

$$p = m \cdot \dot{x}$$

energy

$$H(x, p) = \frac{p^2}{2m} + V(x)$$

Equations of motion

$$\dot{x} = \frac{p}{m} = \frac{\partial H}{\partial p} \qquad \dot{p} = \text{force} = -\frac{\partial V}{\partial x} = -\frac{\partial H}{\partial x}$$

In general, $f(x, p)$ evolves according to

$$\dot{f} = \frac{\partial f}{\partial x} \frac{\partial H}{\partial p} - \frac{\partial H}{\partial x} \frac{\partial f}{\partial p} =: \{f, H\}$$

using Poisson bracket

$$\{-, -\} = \partial_x \wedge \partial_p \qquad \{x, p\} = 1$$

Quantization

Promote to operators on $\mathbb{C}[x]$:

$$x \rightsquigarrow \widehat{x} \cdot \quad p \rightsquigarrow \widehat{p} = -i\hbar\partial_x$$

Canonical commutation relations:

$$[\widehat{x}, \widehat{p}] = \widehat{x}\widehat{p} - \widehat{p}\widehat{x} = i\hbar \quad \{x, p\} = 1$$

Weyl: given $f \in \mathbb{C}[x, p]$, define \widehat{f} by symmetrization:

$$\widehat{(ax + bp)^n} = (a\widehat{x} + b\widehat{p})^n$$

Defines new associative product \star on $\mathbb{C}[x, p]$:

$$\widehat{f \star g} = \widehat{f}\widehat{g} \quad (\mathbb{C}[x, p], \star) \cong \frac{\mathbb{C}\langle \widehat{x}, \widehat{p} \rangle}{\widehat{x}\widehat{p} - \widehat{p}\widehat{x} = i\hbar}$$

The star product

“Explicit” formula:

$$(f \star g)(u) = \int_{v, w \in \mathbb{R}_{x,p}^2} f(v)g(w)e^{iS(u,v,w)/\hbar} dv dw$$

$$S(u, v, w) = 4 \cdot \text{Area} \left(\begin{array}{c} u \\ \triangle \\ v \quad w \end{array} \right)$$

Groenewold 1946, Moyal 1949:

$$f \star g \sim \sum_{n=0}^{\infty} \frac{(i\hbar)^n}{2^n n!} \sum_{i=0}^n (-1)^i \frac{\partial^n f}{\partial x^i \partial p^{n-i}} \frac{\partial^n g}{\partial x^{n-i} \partial p^i}$$

Key point:

$$f \star g - g \star f = i\hbar\{f, g\} + O(\hbar^2)$$

Abstraction

“Phase space” = manifold/variety X equipped with a **Poisson bracket**

$$\{-, -\} : \mathcal{O}_X \times \mathcal{O}_X \rightarrow \mathcal{O}_X$$

Axioms:

- 1 $\{f, g\} = -\{g, f\}$
- 2 $\{f, gh\} = \{f, g\}h + g\{f, h\}$
- 3 $\{f, \{g, h\}\} + \{g, \{h, f\}\} + \{h, \{f, g\}\} = 0$

Examples of Poisson manifolds

- Darboux: $X = \mathbb{R}^{2n}$ with $\{f, g\} = \sum_i \frac{\partial f}{\partial x^i} \frac{\partial g}{\partial p_i} - \frac{\partial g}{\partial x^i} \frac{\partial f}{\partial p_i}$
- (X, ω) symplectic, e.g. $X = T^*Q$.
- Angular momentum: $X = \mathbb{R}_{x,y,z}^3$ with

$$\{x, y\} = z \quad \{y, z\} = x \quad \{z, x\} = y$$

- Linear brackets on $X = \mathbb{R}^n$

$$\{x_i, x_j\} = \sum_k c_{ij}^k x_k \quad \iff \quad \text{Lie algebra } \mathfrak{g} = (\mathbb{R}^n)^*$$

- moduli spaces in gauge theory
- ...

The “deformation quantization” problem

Formulated by Bayen–Flato–Fronsdal–Lichnerowicz–Sternheimer (1978)

A **deformation quantization** of $(X, \{-, -\})$ is a family of associative products \star_{\hbar} such that

$$\begin{aligned}f \star_{\hbar} g &= fg && \text{when } \hbar = 0 \\f \star_{\hbar} g - g \star_{\hbar} f &= \hbar\{f, g\} + O(\hbar^2)\end{aligned}$$

Today: only formal deformations

$$\begin{aligned}\star &: \mathcal{O}_X \times \mathcal{O}_X \rightarrow \mathcal{O}_X[[\hbar]] \\f \star g &= fg + \hbar B_1(f, g) + \hbar^2 B_2(f, g) + \dots\end{aligned}$$

Basic question: Does a quantization always exist?

Answer in symplectic case: yes – Berezin, Deligne, Fedosov, Kirillov, Kostant, Schlichenmaier, Souriau, Toeplitz, ...

Kontsevich formality theorem

Theorem (Kontsevich 1997)


Every smooth Poisson manifold X has a canonical quantization. In fact there is an equivalence

$$\underbrace{\{\text{Poisson brackets on } X\}}_{\sim} \cong \underbrace{\{\text{noncommutative deformations of } \mathcal{O}_X\}}_{\sim}$$

Precise statement is stronger: $CC^*(\mathcal{O}_X) \cong \wedge^\bullet T_X$ as dg Lie algebras

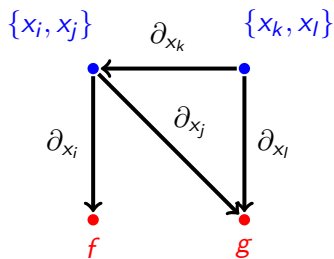
Explicit Feynman expansion when $X = \mathbb{R}^n$:

$$f \star g = fg + \hbar \left(\text{triangle diagram} \right) + \hbar^2 \left(\text{cross diagram} + \text{triangle with arrow} + \text{triangle with loop} \right) + \dots$$


 $=$
(complicated integral) \cdot (derivatives of f, g and $\{-, -\}$)

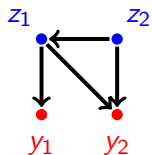
Kontsevich formula: differential operator

Given $\{-, -\}$ in coordinates x_1, \dots, x_n on \mathbb{R}^n , want to compute $f \star g$



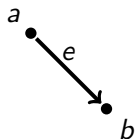
$$\left(\begin{array}{c} \text{derivatives of} \\ f, g \text{ and } \{-, -\} \end{array} \right) := \sum_{i,j,k,l} (\partial_{x_i} f) \cdot (\partial_{x_j} \partial_{x_l} g) \cdot (\partial_{x_k} \{x_i, x_j\}) \cdot \{x_k, x_l\}$$

Kontsevich formula: Feynman integrals



$$\rightsquigarrow \mathfrak{C}_{n,m} = \left\{ \begin{array}{c} \infty \\ \text{circle with points } z_1, z_2, y_1, y_2 \end{array} \right\} / \text{holomorphic iso.}$$

$$\text{e.g. } \mathfrak{C}_{n,2} \cong \left\{ \begin{array}{c} \infty \\ \text{circle with points } 0, 1 \end{array} \right\} \cong \mathbb{H}^n \setminus \{z_i = z_j\}_{i \neq j}$$



$$\rightsquigarrow 2\alpha_e := \frac{d \log(a, \bar{a}; b, \infty)}{2i\pi} + \frac{d \log(a, \bar{a}; \bar{b}, \infty)}{2i\pi}$$

$$\omega := \alpha_{e_1} \wedge \cdots \wedge \alpha_{e_N} \in 2^{-N} \mathbb{Z} \left\langle \frac{d \log f}{2i\pi} \mid f \text{ a cross ratio} \right\rangle \subset \Omega^\bullet(\mathfrak{C}_{n,m})$$

$$\left(\begin{array}{c} \text{complicated} \\ \text{integral} \end{array} \right) := \int_{\mathfrak{C}_{n,m}} \omega \in \mathbb{R}$$

Recovering Groenewold–Moyal

$$\{f, g\} = \frac{\partial f}{\partial x} \frac{\partial g}{\partial p} - \frac{\partial g}{\partial x} \frac{\partial f}{\partial p} \quad \left(\begin{array}{cc} \{x, x\} & \{x, p\} \\ \{p, x\} & \{p, p\} \end{array} \right) = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\begin{aligned} f \star g &= fg + \hbar \left(\text{Diagram 1} \right) + \hbar^2 \left(\text{Diagram 2} + \text{Diagram 3} + \text{Diagram 4} \right) + \dots \\ &= fg + \hbar \left(\text{Diagram 1} \right) + \hbar^2 \left(\text{Diagram 2} \right) + \hbar^3 \left(\text{Diagram 3} \right) + \dots \\ &= \sum_{n=0}^{\infty} \frac{(i\hbar)^n}{2^n n!} \sum_{i=0}^n (-1)^i \frac{\partial^n f}{\partial x^i \partial p^{n-i}} \frac{\partial^n g}{\partial x^{n-i} \partial p^i} \end{aligned}$$

Diagram 1: A blue dot at the top with two arrows pointing down to two red dots.

 Diagram 2: Two blue dots at the top with four arrows pointing down to two red dots, crossing each other.

 Diagram 3: Two blue dots at the top with four arrows pointing down to two red dots, crossing each other.

 Diagram 4: Two blue dots at the top with four arrows pointing down to two red dots, crossing each other. Dashed arrows indicate a loop between the two blue dots.

Linear case

$$\{x_i, x_j\} = \sum c_{ij}^k x_k \quad \leftrightarrow \quad \text{Lie algebra } \mathfrak{g}$$

Similar analysis:

- Series truncates for $f, g \in \mathbb{C}[x_i]$
- Can compute

$$x_i \star x_j - x_j \star x_i = \hbar \sum c_{ij}^k x_k$$

- Conclude

$$(\mathbb{C}[x_i], \star_{\hbar}) \cong \frac{\mathbb{C}\langle x_i \rangle}{x_i x_j - x_j x_i = \hbar \sum c_{ij}^k x_k} =: U(\mathfrak{g}, \hbar)$$

Quadratic case

$$\{X, P\} = XP$$

$$X \star P = q(\hbar)XP \quad P \star X = q(-\hbar)XP$$

Our software:

$$q(\hbar) = 1 + \frac{\hbar}{2} + \frac{\hbar^2}{24} - \frac{\hbar^3}{48} - \frac{\hbar^4}{1440} + \frac{\hbar^5}{480} + \left(\frac{251\zeta(3)^2}{2048\pi^6} - \frac{17}{184320} \right) \hbar^6 + \dots$$

Nevertheless: algebra determined by

$$X \star P = \frac{q(\hbar)}{q(-\hbar)} P \star X = e^{\hbar} P \star X$$

Morally: $X = e^x$ and $P = e^p$ where $\{x, p\} = 1$.

Special values of Riemann zeta

$$\zeta(s) = \sum_{k \geq 1} \frac{1}{k^s}$$

Theorem (Euler 1735): $\zeta(2m) = (-1)^{m+1} \frac{B_{2m}(2\pi)^{2m}}{2(2m)!} \in \mathbb{Q}\pi^{2m}$

Open Question: Is $\zeta(2m+1) \in \mathbb{Q}(\pi)$?

Conjecture: $\pi, \zeta(3), \zeta(5), \zeta(7), \dots$ are algebraically independent over \mathbb{Q} .

Theorem (Apéry 1978): $\zeta(3) \notin \mathbb{Q}$

Theorem ((Ball-)Rivoal 2000): Infinitely many $\zeta(3), \zeta(5), \zeta(7), \dots \notin \mathbb{Q}$

Theorem (Zudilin 2000): At least one of $\zeta(5), \zeta(7), \zeta(9), \zeta(11) \notin \mathbb{Q}$

MZVs (Euler, Écalle, Zagier, ...)

Definition

A **normalized multiple zeta value (MZV)** of **weight n** is a number of the form

$$\tilde{\zeta}(n_1, \dots, n_d) = \frac{1}{(2i\pi)^n} \sum_{k_1 > k_2 > \dots > k_d \geq 1} \frac{1}{k_1^{n_1} k_2^{n_2} \dots k_d^{n_d}} \in \begin{cases} \mathbb{R} & n \text{ even} \\ i\mathbb{R} & n \text{ odd} \end{cases}$$

where $n_1 \geq 2$ and $n_1 + \dots + n_d = n$.

Additional “honourary” normalized MZVs:

- $1 \in \mathbb{R}$ has weight 0
- $\frac{1}{2} = \frac{i\pi}{2i\pi} \in \mathbb{R}$ has weight 1

Algebra of MZVs

$$\tilde{\mathcal{Z}} := \mathbb{Z} \cdot \{\text{normalized MZVs}\} \subset \mathbb{C}$$

Weight filtration:

$$\begin{array}{ccccccc} \tilde{\mathcal{Z}}_0 & \subset & \tilde{\mathcal{Z}}_1 & \subset & \tilde{\mathcal{Z}}_2 & \subset & \dots \subset \tilde{\mathcal{Z}} \\ \parallel & & \parallel & & \parallel & & \\ \mathbb{Z} & \subset & \mathbb{Z} \cdot \frac{1}{2} & \subset & \mathbb{Z} \cdot \underbrace{\frac{\zeta(2)}{(2i\pi)^2}}_{=\frac{-1}{24}} & \subset & \dots \end{array}$$

Shuffle product:

$$\tilde{\mathcal{Z}}_m \tilde{\mathcal{Z}}_n \subset \tilde{\mathcal{Z}}_{m+n}$$

e.g.

$$\tilde{\zeta}(m)\tilde{\zeta}(n) = \tilde{\zeta}(m, n) + \tilde{\zeta}(n, m) + \tilde{\zeta}(n + m)$$

How many MZVs are there?

For unnormalized MZVs:

- \mathbb{Q} -dimension of weight spaces conjectured by Zagier
 - ▶ Proven to be an upper bound (Terasoma, Deligne–Goncharov)
- \mathbb{Q} -basis conjectured by Hoffman: $\zeta(2s$ and $3s)$
 - ▶ Proven to generate (Brown)

For normalized MZVs:

\mathbb{Z} -module generators of $\tilde{\mathcal{Z}}_n$							
n	0	1	2	3	4	5	6
real	1	$\frac{1}{2}$	$\frac{1}{24}$	$\frac{1}{48}$	$\frac{1}{5760}$	$\frac{1}{11520}$	$\frac{1}{2903040}$
							$\frac{\zeta(3)^2}{128\pi^6}$
imaginary				$\frac{i\zeta(3)}{8\pi^3}$	$\frac{i\zeta(3)}{16\pi^3}$	$\frac{i\zeta(3)}{192\pi^3}$	$\frac{i\zeta(3)}{384\pi^3}$
						$\frac{i\zeta(5)}{64\pi^5}$	$\frac{i\zeta(5)}{128\pi^5}$

Ubiquity of MZVs

- **Quantum groups:** coefficients of Drinfel'd associator
- **Knot theory:** coefficients of Kontsevich integral
- **Homotopical algebra:** formality of the operad E_2
- **Algebraic geometry:** periods integrals on moduli space $\mathcal{M}_{0,N}$
(Brown 2006, conj. by Goncharov–Manin)
- **Physics:** values of certain Feynman integrals
- ...

Theorem (Brown 2011, building on Deligne–Goncharov, Levine, Voevodsky, Zagier, ...)

All periods of unramified mixed Tate motives lie in $\mathbb{Q}\tilde{\mathcal{Z}}[\frac{1}{2i\pi}]$.

Main result

$$\mathfrak{C}_{n,m} = \left\{ \text{circle with } n \text{ points} \right\} / \text{holomorphic iso.}$$

$$\mathcal{A}^\bullet(\mathfrak{C}_{n,m}) := \mathbb{Z} \left\langle \frac{d \log f}{2i\pi} \mid f \text{ a cross ratio} \right\rangle \subset \Omega^\bullet(\mathfrak{C}_{n,m})$$

Theorem (Banks–Panzer–P.)

Suppose that $\omega \in \mathcal{A}^\bullet(\mathfrak{C}_{n,m})$ is absolutely integrable. Then

$$\int_{\mathfrak{C}_{n,m}} \omega \in \begin{cases} \tilde{\mathcal{Z}}_{n+m-2} & m > 0 \\ \tilde{\mathcal{Z}}_{n-1} & m = 0 \end{cases}$$

Corollary (case $m = 2$)

Coefficients at order \hbar^n in Kontsevich's star product lie in $4^{-n} \tilde{\mathcal{Z}}_n \cap \mathbb{R}$

Alternate definitions of MZVs

$$\tilde{\zeta}(n_1, \dots, n_d) = \frac{1}{(2i\pi)^n} \sum_{k_1 > k_2 > \dots > k_d \geq 1} \frac{1}{k_1^{n_1} k_2^{n_2} \dots k_d^{n_d}} = L_{n_1, \dots, n_d}(1)$$

in terms of **multiple polylogarithm**

$$L_{n_1, \dots, n_d}(z) := \frac{1}{(2i\pi)^n} \sum_{k_1 > k_2 > \dots > k_d \geq 1} \frac{z^{k_1}}{k_1^{n_1} k_2^{n_2} \dots k_d^{n_d}}$$

e.g.

$$L_1(z) = \sum_{k \geq 1} \frac{z^k}{k} = \frac{\log(1-z)}{2i\pi} \quad L_2(z) = \text{dilogarithm}$$

Alternate notation:

$$n_1, \dots, n_d \quad \leftrightarrow \quad s_1 \dots s_n = \underbrace{00 \dots 01}_{n_1} \underbrace{00 \dots 01}_{n_2} \dots \underbrace{00 \dots 01}_{n_d}$$

Check:

$$dL_{s_1 \dots s_n} = (-1)^{s_1} \frac{L_{s_2 \dots s_n} dz}{2i\pi(z - s_1)}$$

Alternate definitions of MZVs, II

Rewrite

$$\tilde{\zeta}(n_1, \dots, n_d) = L_{s_1 \dots s_n}(1) \quad dL_{s_1 \dots s_n} = (-1)^{s_1} \frac{L_{s_2 \dots s_n} dz}{2i\pi(z - s_1)}$$

and therefore (Kontsevich, Le-Murakami)

$$\tilde{\zeta}(n_1, \dots, n_d) = (-1)^d \underbrace{\int_0^1 \frac{dt_1}{2i\pi(t_1 - s_1)} \int_0^{t_1} \frac{dt_2}{2i\pi(t_2 - s_2)} \cdots \int_0^{t_{n-1}} \frac{dt_n}{2i\pi(t_n - s_n)}}_{\int_0^1 s_1 \cdots s_n}$$

Chen iterated integral

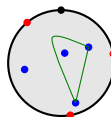
NB: diverges if $s_1 = 1$ or $s_n = 0$, so “regularize”: $\log(\epsilon) = 0$

Polylogs on $\mathfrak{C}_{m,n}$ (BPP), cf. Brown, Goncharov for $\mathfrak{M}_{0,N}$

$$\mathfrak{C}_{n,m} = \left\{ \text{disk with } m \text{ red points and } n \text{ blue points} \right\} / \text{holomorphic iso.}$$

Choose $s_0, s_1, \dots, s_{n+1} \in \{z_i, \bar{z}_i, y_i\}$, define “disk polylog” (multivalued!)

$$L_{s_0; s_1 \cdots s_n; s_{n+1}} : \mathfrak{C}_{n,m} \rightarrow \mathbb{C}$$



$$\mapsto \int_{s_0}^{s_{n+1}} s_1 \cdots s_n$$

regularizing divergences via Deligne's tangential base points.

These functions and their differentials generate a locally constant subsheaf

$$\mathcal{A}^\bullet(\mathfrak{C}_{n,m}) \subset \mathcal{U}_{\mathfrak{C}_{n,m}}^\bullet \subset \Omega_{\mathfrak{C}_{n,m}}^\bullet$$

with monodromy unipotent for the weight filtration.

$$\text{Constants: } \tilde{\mathcal{Z}} \subset \mathcal{U}_{\mathfrak{C}_{n,m}}^0$$

de Rham isomorphism

Theorem (BPP “de Rham theorem for disk polylogs”)

\mathcal{U}^\bullet is a resolution of the constant sheaf $\tilde{\mathcal{Z}}$ by acyclic local systems. Hence

$$H^\bullet(\mathcal{U}^\bullet(\mathfrak{C}_{n,m}), d) \cong H^\bullet(\mathfrak{C}_{n,m}; \tilde{\mathcal{Z}}).$$

Sketch of proof.

Induction on n, m via $f : \mathfrak{C}_{n,m} \rightarrow \mathfrak{C}_{n-k, m-j}$.

Resolution: \mathbb{Z} -linear lift of Brown’s Poincaré lemma via fibrewise KZ equation $dL = L' \cdot dz / (z - s)$.

Acyclic: have $\mathfrak{C}_{n,m} = K(\text{PureBraids}_n, 1)$, show group cohomology of the monodromy representation vanishes (i.e. $R^{>0} f_* \mathcal{U}_{\mathfrak{C}_{n,m}}^\bullet = 0$). □

Corollary

Every volume form in $\mathcal{U}^\bullet(\mathfrak{C}_{n,m})$ has a primitive in $\mathcal{U}^\bullet(\mathfrak{C}_{n,m})$.

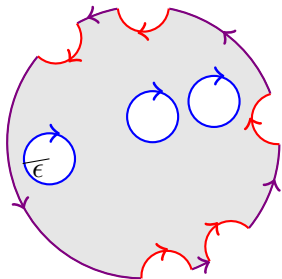
Integration

Theorem (BPP “Fubini theorem for disk polylogs”)

Given $f : \mathfrak{C}_{n,m} \rightarrow \mathfrak{C}_{n-k,m-j}$ and integrable $\omega \in \mathcal{U}^\bullet(\mathfrak{C}_{n,m})$, have

$$\int_{\mathfrak{C}_{n,m}} \omega = \int_{\mathfrak{C}_{n-k,m-j}} \left(\int_{\text{fibres}} \omega \right) \quad \int_{\text{fibres}} \omega \in \mathcal{U}^\bullet(\mathfrak{C}_{n-k,m-j})$$

and weight drops by k . Main theorem: $\omega \in \mathcal{A}^\bullet(\mathfrak{C}_{n,m})$ and $f : \mathfrak{C}_{n,m} \rightarrow \text{pt.}$



$$\begin{aligned} & \int_{\text{disk}} \frac{Ldz \wedge d\bar{z}}{(z-s)(\bar{z}-z)} \\ &= \lim_{\epsilon \rightarrow 0} \int_{\partial_\epsilon \text{disk}} \tilde{L} \frac{dz}{z-s} \\ &= \sum \text{Res} + \int_{\text{outer cycle}} \end{aligned}$$

estimates + unipotent monodromy
 \rightsquigarrow weight drop

Motivic directions

Conjecture (BPP)

Coefficients of the star product at \hbar^n generate $\tilde{\mathcal{Z}}_n$.

Strategy: operadic motivic lift (in progress with Dupont and Panzer)

Convergence of power series? Motivic Galois action?

Thank you!